

# Resonance crossing and emittance growth

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KEK

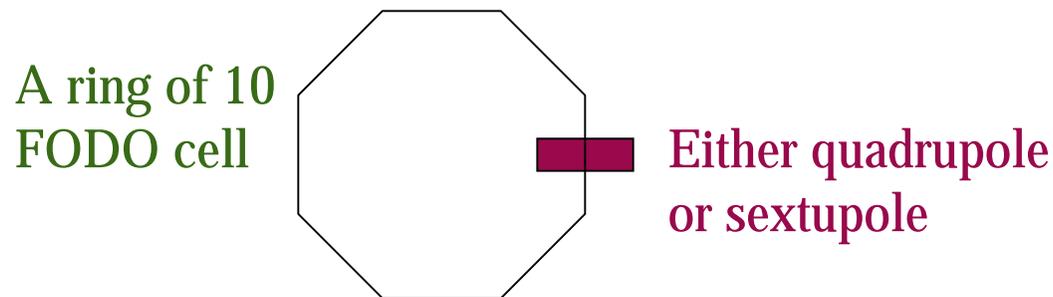
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# Purpose

- In non-scaling FFAG, there is wide excursion of tune from injection to extraction. A beam crosses major resonance several times.
- In scaling FFAG, zero chromaticity condition is not perfect. Resonance crossing occurs during acceleration.
- We want to know
  - how a beam is affected by resonances crossing.
  - if there is any necessary condition to avoid the effects.

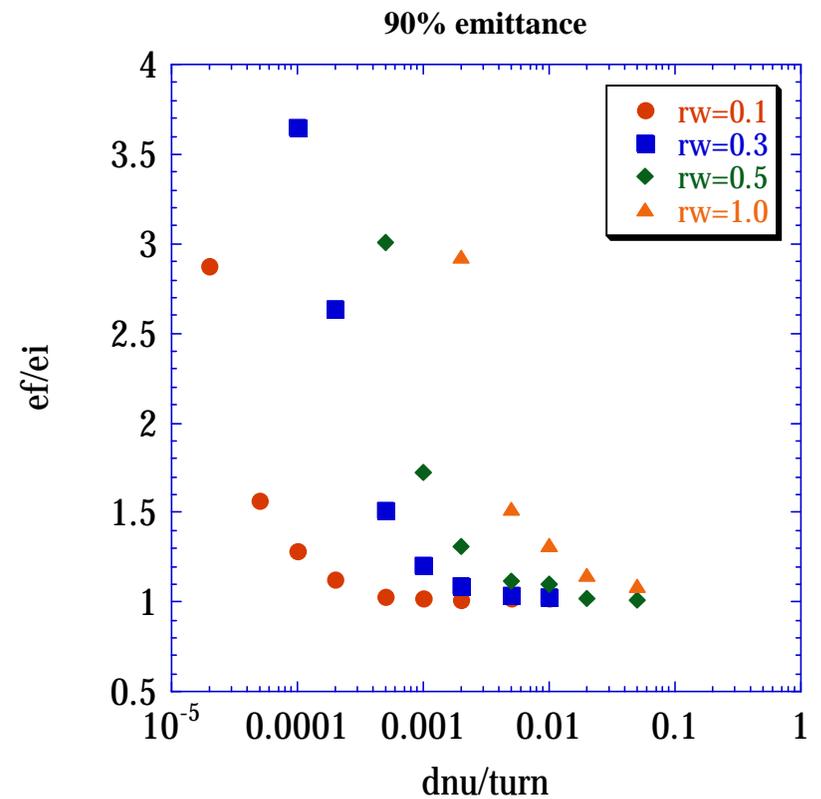
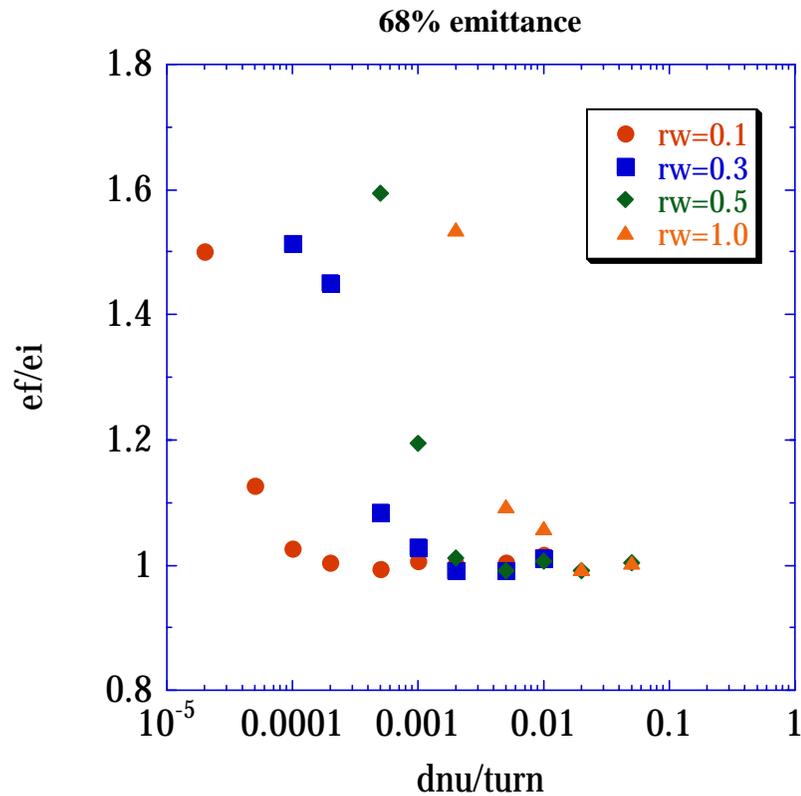
# Results presented at BNL workshop

- We study resonance crossing (1-D) on a simplified model.
  - 10 FODO cell and single source of resonance.

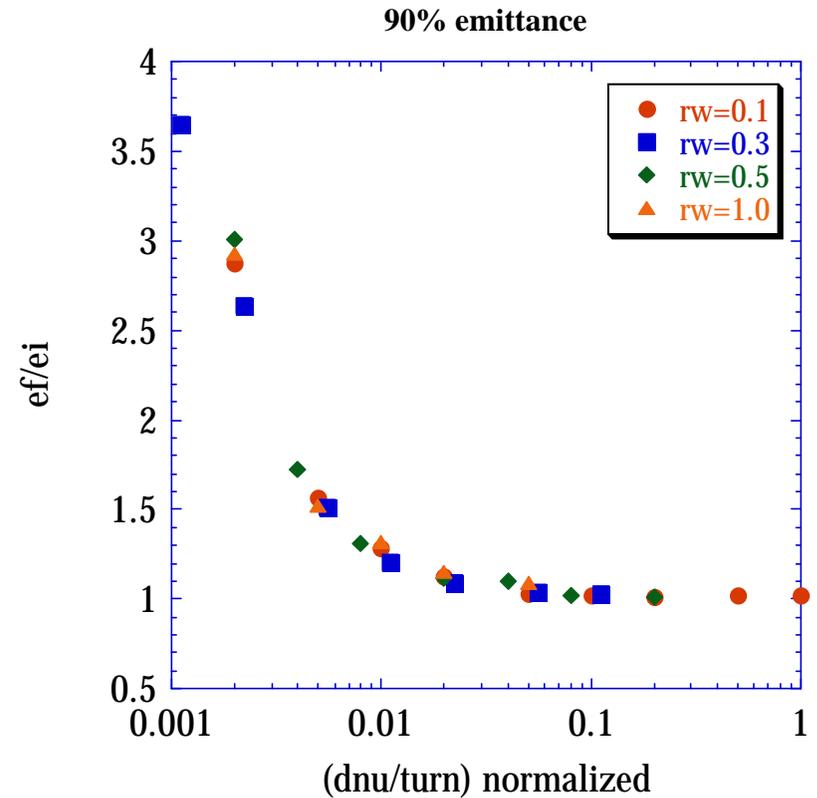
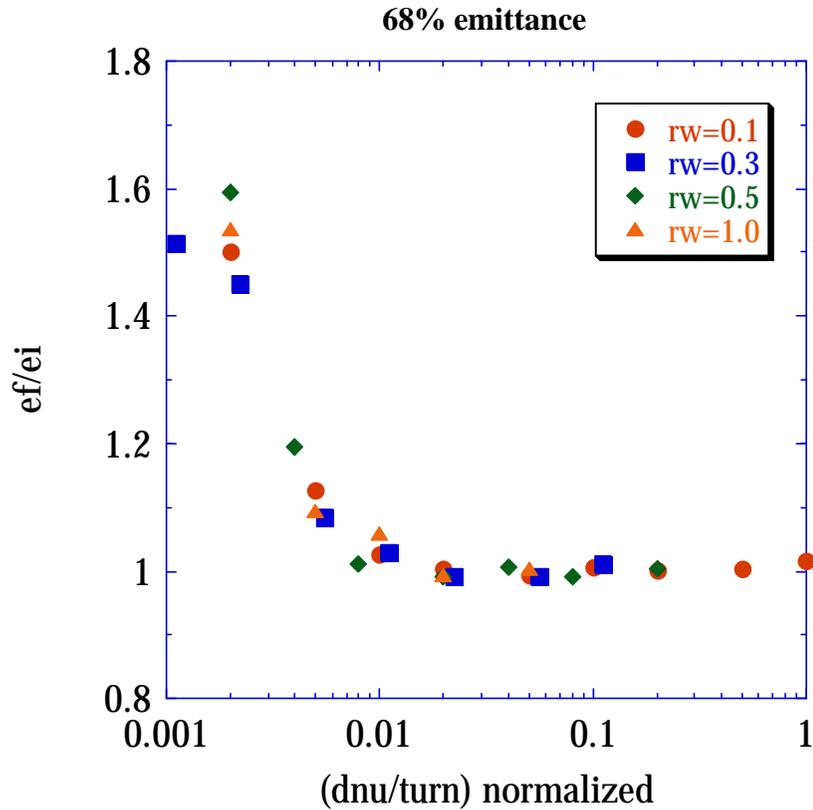


- At half integer resonance crossing,
  - Emittance growth is the same as long as  $(\text{Tune shift per turn})/(\text{Resonance width})$  is the same.
  - To avoid any effects by crossing,  $(\text{Tune shift per turn}) > (\text{Resonance width})$ . Reasonable!
- At third (and fourth) integer resonance crossing,
  - Emittance growth can be parameterized with  $\text{Tune shift per turn}/(\text{sextupole strength})^2$ . Why so?

# Simulation results (at the BNL workshop)



# After normalization with $k^2$ <sup>2</sup> (at the BNL workshop)



# Movie of phase space evolution

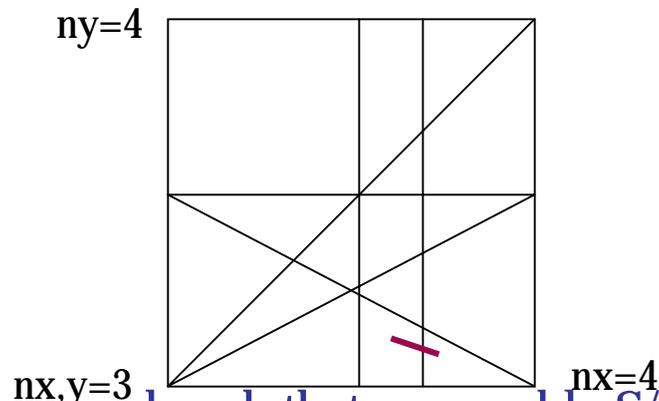
- Single resonance example

# Recent progress

- Take HIMAC lattice as an example, simulation and experiment have been done to see the following parameter dependence when a beam cross 3rd order resonance.
  - Resonance strength (sextupole strength)
  - Crossing speed
- The following quantities have been measured.
  - Beam loss (simulation and experiment)
  - Emittance (simulation)
  - Beam profile (experiment)

# Detailed procedure

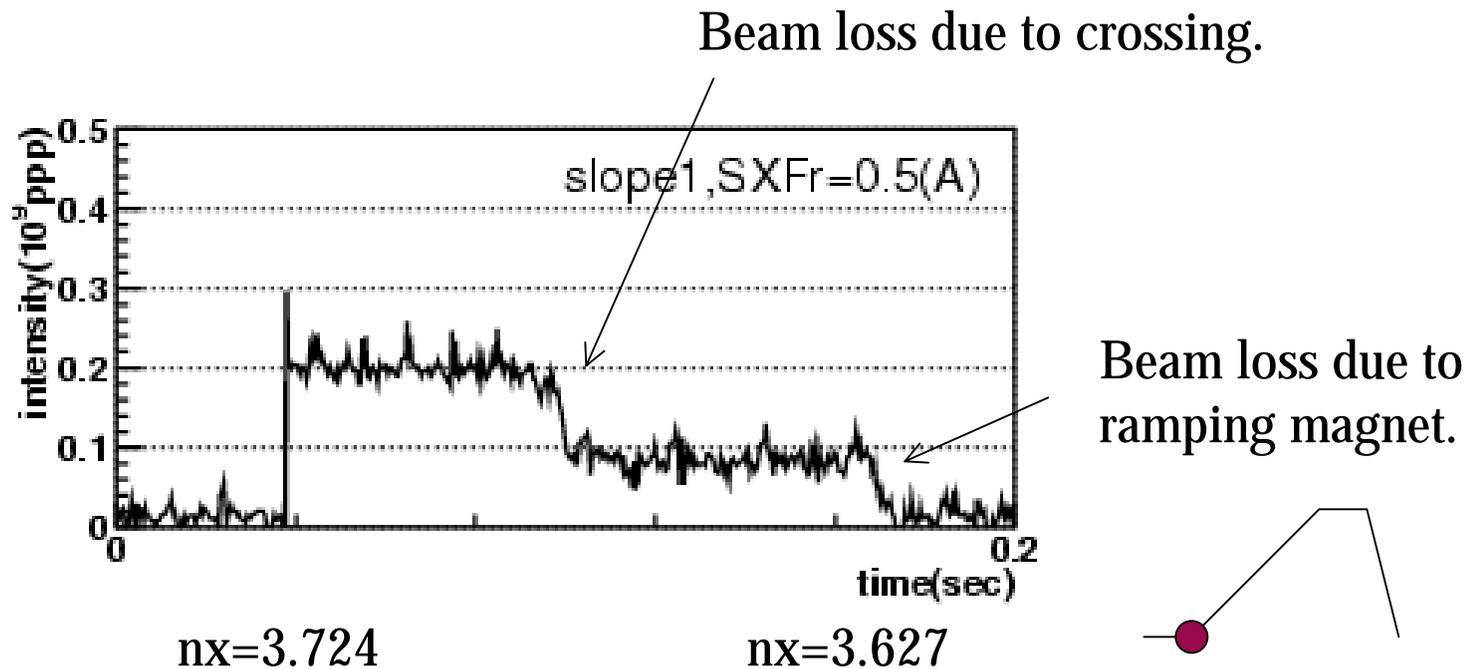
- HIMAC lattice has 6 fold symmetry. Excite one or two sextupole magnets which create 11th harmonic of resonance at  $n_x=3.666$  ( $3n_x=11$ ).
  - cf. That sextupole is used to excite  $n_x=3.666$  for slow extraction.
- Bare tune is varied linearly after injection with QF.  $n_x=3.666$  is crossed in a middle way.



- Beam intensity is lowered such that reasonable S/N ratio is still obtained.
- In simulation, emittance growth and beam loss are observed. In experiment, beam loss and beam profile are measured.

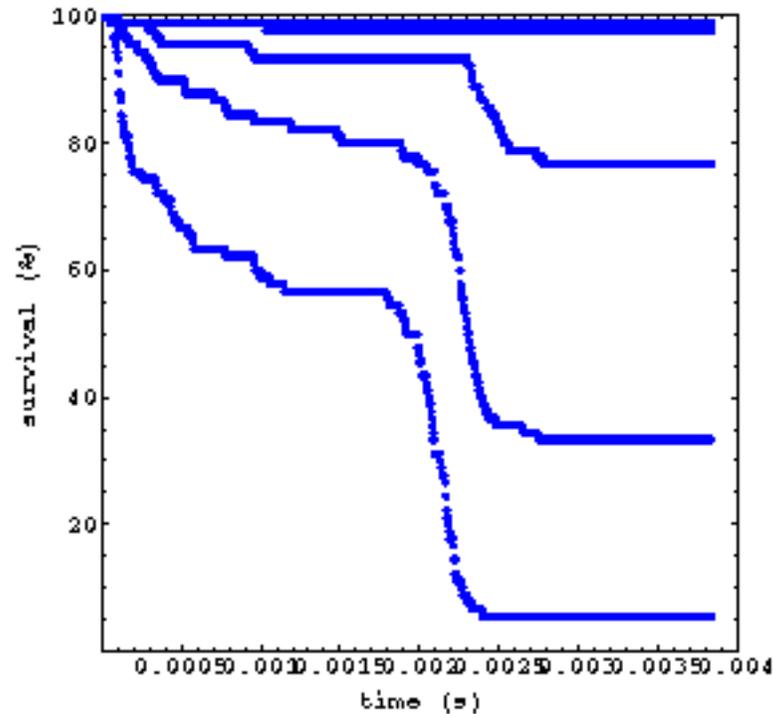
# Beam loss in HIMAC beam study

Beam intensity is lowered to avoid tune spread while keeping a good S/N ratio.



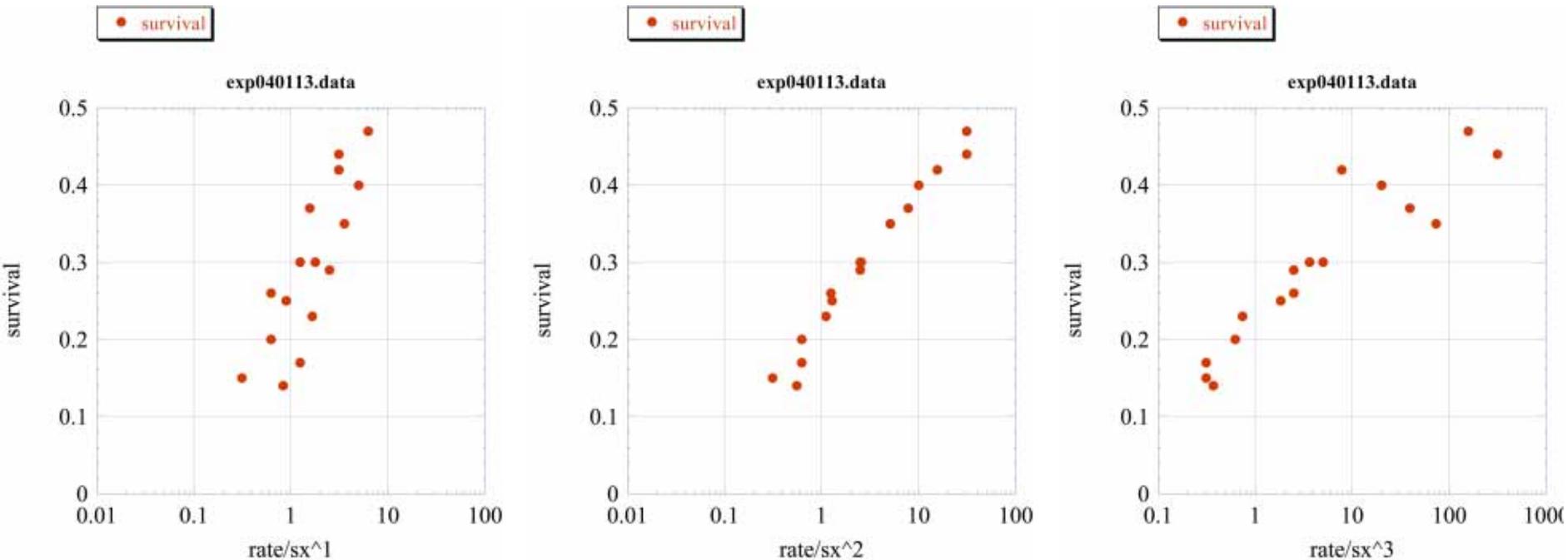
# Beam loss due to resonance crossing

- Simulation results of resonance crossing and beam loss. Sextupole strength is a parameter. ( $K_2=0.01, 0.02, 0.05, 0.1, 0.2$ )



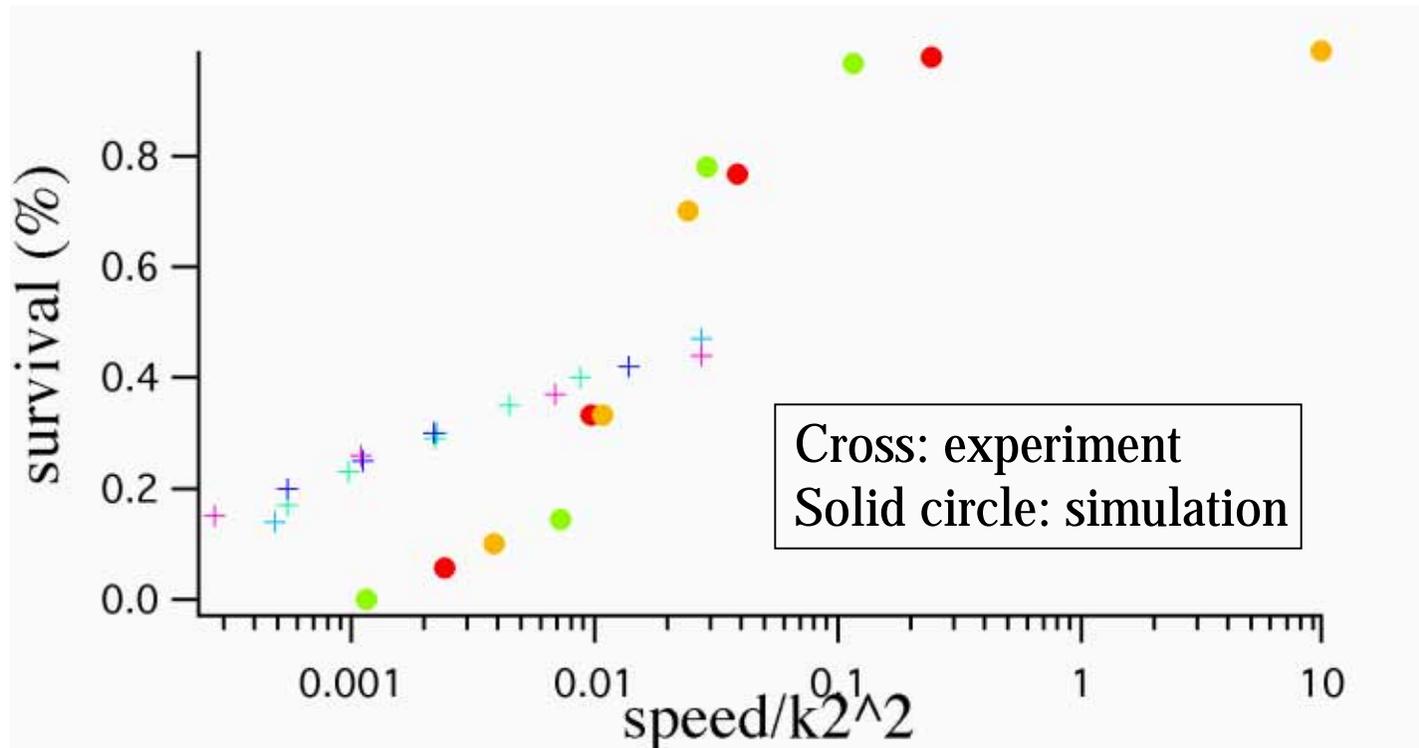
# Best fit of experimental results

- Three ways of parameterization. From left to right, x-axis is  $\text{Speed}/k^2^1$ ,  $\text{Speed}/k^2^2$ ,  $\text{Speed}/k^2^3$ .
- When we choose  $\text{Speed}/k^2^2$  as x-axis, all experimental results sit on the same line.



# Comparison between experiment and simulation

Both simulation and experimental data can be parameterized with  $\text{speed}/k^2$ . However, absolute survival ratio does not agree.



# Summary

- Beam loss due to 3rd order resonance crossing can be normalized with  $\text{speed}/k_2^2$ , where speed is a speed of crossing and  $k_2$  is strength of sextupole.
- Although both simulation and experimental results show  $\text{speed}/k_2^2$  is the best parameterization, they do not agree quantitatively.
  - Initial emittance is not well known in experiment.
  - Other resonance sources in the real HIMAC lattice?

## Further study

- In scaling FFAG, there are many other nonlinear resonances. It may not be a right picture to pick up single resonance.
- For example, when there is amplitude dependent tune shift, three fixed points in addition to the origin will appear. Particle trapping by those islands has to be considered as a source of beam loss.
- Simulation and experiment on the trapping by resonance islands are also studied (as a theses subject.)